

COMPITO DEL 20-12-11

Esercizio 1: Gas ideale di Fermioni in 1D a $T=0$

① L'Hamiltoniana di singola particella è

$$h^{(1)} = -\frac{\hbar^2}{2m} \frac{d^2}{dx^2}$$

e la parte orbitale della funzione d'onda di singola particella soddisfa

$$-\frac{\hbar^2}{2m} \varphi''(x) = E \varphi(x) \Rightarrow \varphi''(x) + k^2 \varphi(x) = 0$$
$$k^2 = \frac{2m}{\hbar^2} E$$

$$\varphi(x) = A \sin kx + B \cos kx \quad \varphi(0) = 0 \Rightarrow B = 0$$

$$\varphi(L) = 0 \Rightarrow k = \frac{\pi}{L} n$$

• orbitali $\left[\phi_m(x, s_z) = \sqrt{\frac{2}{L}} \sin\left(\frac{\pi}{L} n x\right) \chi_{s_z} \right]_{n \in \mathbb{N}}$
 $n \in \mathbb{N} \quad (n = 1, 2, 3, \dots)$

• energie

$$\left[\bar{E}_n = \frac{\hbar^2}{2m} \left(\frac{\pi}{L}\right)^2 n^2 \right]$$

② I k_n sono spaziali di π/L e le energie $\propto k_n^2 \propto n^2$; occupando ciascun k_n con 2 elettroni ($s_z = \pm \hbar/2$) a partire dal più piccolo ($n=1$) otteniamo

$$N = \frac{2K_F}{\pi/L} \Rightarrow \left[K_F = \frac{N}{2} \frac{\pi}{L} = M \frac{\pi}{L} \right]$$

$$\left[E_F = \frac{\hbar^2}{2m} \left(\frac{\pi}{L} \right)^2 M^2 \right] = \frac{\hbar^2}{2m} \pi^2 \left(\frac{\tilde{n}}{2} \right)^2 \quad \tilde{n} = \frac{N}{L}$$

③
$$E = 2 \frac{\hbar^2}{2m} \frac{\pi^2}{L^2} \sum_{n=1}^M n^2 = \frac{\hbar^2}{2m} \frac{\pi^2}{L^2} \frac{1}{3} M(M+1)(2M+1)$$

$$E = \frac{\hbar^2}{24m} \frac{\pi^2}{L^2} N(N+2)(N+1)$$

$$\left[\frac{E}{N} = \frac{\hbar^2}{24m} \frac{\pi^2}{L^2} (N+1)(N+2) \right]$$

④
$$\mu = \frac{\partial E}{\partial N} = \frac{\hbar^2}{24m} \frac{\pi^2}{L^2} \frac{\partial}{\partial N} [N^3 + 3N^2 + 2N]$$

$$\mu = \frac{\hbar^2}{24m} \frac{\pi^2}{L^2} [3N^2 + 6N + 2] = \frac{\hbar^2}{2m} \frac{\pi^2}{L^2} \left(\frac{N}{2} \right)^2 \left[1 + \frac{2}{N} + \frac{2}{3N^2} \right]$$

$$\left[\mu = E_F \left[1 + \frac{2}{N} + \frac{2}{3N^2} \right] \right]$$

$$\left[\lim_{N \rightarrow \infty} \mu(N) = E_F \quad \frac{N}{L} = \tilde{n} = \text{cost.} \right]$$

Esercizio 2: Fonnini alla Debye in un cristallo 1D.

$$\begin{aligned} \textcircled{1} \quad Q &= \sum_{\{n_k\}} \exp\left[-\beta \sum_k \hbar \omega_k n_k\right] \\ &= \prod_k \sum_{n_k=0}^{\infty} \exp\left[-\beta \hbar \omega_k n_k\right] \\ &= \prod_k \sum_{n_k=0}^{\infty} \left(\exp\left[-\beta \hbar \omega_k\right]\right)^{n_k} \end{aligned}$$

$$\boxed{Q = \prod_k \frac{1}{1 - e^{-\beta \hbar \omega_k}}}$$

$$\textcircled{2} \quad E = \frac{\partial(\beta A)}{\partial \beta} = \frac{\partial}{\partial \beta} (-\ln Q) = \frac{\partial}{\partial \beta} \sum_k \ln(1 - e^{-\beta \hbar \omega_k})$$

$$E = \sum_k \frac{\hbar \omega_k e^{-\beta \hbar \omega_k}}{1 - e^{-\beta \hbar \omega_k}} = \sum_k \frac{\hbar \omega_k}{e^{\beta \hbar \omega_k} - 1}$$

For $L \rightarrow \infty$ (P.B.C.)

$$\frac{E}{L} = \frac{1}{L} \int_{-k_D}^{k_D} \frac{d\alpha}{\frac{2\pi}{L}} \frac{\hbar c |\alpha|}{e^{\beta \hbar c |\alpha|} - 1}$$

$$= \frac{2}{2\pi} \int_0^{k_D} d\alpha \frac{\hbar c \alpha}{e^{\beta \hbar c \alpha} - 1}$$

$$\boxed{t = \beta \hbar c \alpha}$$

$$\frac{E}{L} = \frac{1}{\pi} \frac{(k_B T)^2}{\hbar c} \int_0^{T_0/T} dt \frac{t}{e^t - 1}$$

$$\frac{\hbar c k_D}{k_B T} = \frac{T_0}{T}$$

$$\lim_{T \rightarrow 0} \frac{E}{L} \approx \frac{1}{\pi} \frac{(k_B T)^2}{\hbar c} \int_0^{\infty} dt \frac{t}{e^t - 1} = \frac{\pi^2}{6} \frac{1}{\pi} \frac{(k_B T)^2}{\hbar c}$$

• $T \rightarrow 0$

$$\boxed{\frac{E}{L} = \frac{\pi}{6} \frac{(k_B T)^2}{\hbar c}}$$

$$\textcircled{3} \quad C_V = \frac{\partial(E/L)}{\partial T_1} = \frac{\partial}{\partial T_1} \left(\frac{\pi}{6} \frac{(k_B T_1)^2}{\hbar c} \right) = \frac{\pi}{3} \frac{k_B^2 T_1}{\hbar c}$$

$$\boxed{C_V = \frac{\pi}{3} \frac{k_B^2 T}{\hbar c} k_B}$$

$\textcircled{4}$ Ricordando $k_0 \leq k \leq k_D$ e $\Delta k = 2\pi/L$

$$N = \frac{2 k_D}{2\pi/L} \Rightarrow k_D = \frac{\pi}{L} N = \pi n$$

$$\boxed{k_D = \pi n}$$

$$k_D = \frac{\pi}{100 \text{ \AA}} = \frac{\pi}{10^{-6} \text{ cm}} = 3.14 \times 10^6 \text{ cm}^{-1}$$

$$\boxed{k_D = 3.14 \times 10^6 \text{ cm}^{-1}}$$