

$$\textcircled{1} g(E) = \frac{2}{L} \sum_p \delta(E - \alpha|p|), \quad -\infty < p < \infty$$

$$g(E) = \frac{2}{L} \int_{-\infty}^{\infty} \frac{dp}{h} \delta(E - \alpha|p|) = \frac{2}{h} 2 \int_0^{\infty} dE \delta(E - \alpha p) \quad \boxed{E = \alpha p}$$

$$= \frac{4}{h\alpha} \int_0^{\infty} dt \delta(E - t) = \frac{4}{h\alpha} \Theta(E)$$

$$\textcircled{2} \rho = \int dE g(E) \frac{1}{\frac{e^{\beta E}}{z} + 1} = \frac{4}{h\alpha} \int_0^{\infty} dE \frac{1}{\frac{e^{\beta E}}{z} + 1} = \frac{4k_B T}{h\alpha} \int_0^{\infty} dy \frac{z e^{-y}}{1 + z e^{-y}}$$

$$= \frac{1}{l} \int_0^{\infty} dy \frac{\partial}{\partial y} [-\ln(1 + z e^{-y})] = \frac{1}{l} \ln(1 + z), \quad l = \frac{h\alpha}{4k_B T}$$

$$\boxed{\rho = \frac{1}{l} \ln(1 + z)}$$

$$\textcircled{3} z = e^{\rho l} - 1 = e^{\beta \mu}$$

$$\beta \mu = \ln(e^{\rho l} - 1)$$

$$\mu = k_B T \ln[e^{\rho l} (1 - e^{-\rho l})] = k_B T \rho l + k_B T \ln(1 - e^{-\rho l})$$

$$= \rho \frac{h\alpha}{4} + k_B T \ln(1 - e^{-\frac{\rho h\alpha}{4k_B T}})$$

$$\lim_{T \rightarrow 0} \mu(T, \rho) = \rho \frac{h\alpha}{4} \equiv \mu_0(\rho)$$

$$\textcircled{4} \quad N = 2 \sum_{|p| < p_F} = 2 \int_{-p_F}^{p_F} \frac{dp}{\frac{h}{L}} = \frac{2L}{h} 2 \int_0^{p_F} dp = \frac{4L}{h} p_F$$

$$p_F = \frac{N}{L} \frac{h}{4} \quad \epsilon_F = \frac{p_F^2}{4}$$

$$\mu_0(p) = \epsilon_F(p)$$

ESERCIZIO 2

$$\begin{aligned}
 \textcircled{1} \quad g(E) &= \frac{1}{V} \sum_{\vec{p}} \delta(E - \frac{p^2}{2m}) = \frac{1}{V} \int \frac{d^3p}{(\frac{h}{L})^3} \delta(E - \frac{p^2}{2m}) \\
 &= \frac{1}{h^3} \int d\Omega \int_0^\infty dp p^2 \delta(E - \frac{p^2}{2m}) = \frac{4\pi}{h^3} (2m)^{3/2} \int_0^\infty d\sqrt{s} s^2 \delta(E - s) \\
 &= \frac{8\pi^2}{3h^3} (2m)^{3/2} \int_0^E ds \frac{s^{3/2}}{2} \delta(E - s) = \frac{4\pi^2}{3h^3} (2m)^{3/2} E^{3/2} \Theta(E) \equiv C E^{3/2} \Theta(E)
 \end{aligned}$$

$$\textcircled{2} \quad \rho = \langle \frac{N}{V} \rangle = \frac{1}{V} \frac{Z}{1-Z} + \int dE C E^{3/2} \frac{1}{\frac{e^{\beta E}}{Z} - 1} \equiv \rho_0 + \rho_1$$

$$\rho_1 = C (k_B T)^{5/2} \int_0^\infty dy \frac{y^{3/2}}{\frac{e^y}{Z} - 1} = C (k_B T)^{5/2} \int_0^\infty dy y^{3/2} \frac{Z e^{-y}}{1 - Z e^{-y}} =$$

$$C (k_B T)^{5/2} \sum_{n=0}^\infty (Z)^{n+1} \int_0^\infty dy y^{3/2} e^{-(n+1)y} = C (k_B T)^{5/2} \sum_{n=1}^\infty (Z)^{n-1} Z^n \int_0^\infty dy y^{3/2} e^{-ny}$$

$$= C (k_B T)^{5/2} \sum_{n=1}^\infty (Z)^{n-1} \frac{Z^n}{n^{5/2}} \int_0^\infty ds s^{3/2} e^{-s} = (k_B T)^{5/2} C \Gamma(\frac{5}{2}) \sum_{n=1}^\infty (Z)^{n-1} \frac{Z^n}{n^{5/2}}$$

$$= \frac{4\pi^2}{3h^3} (2m)^{3/2} (k_B T)^{5/2} \frac{3}{2} \frac{1}{2} \sqrt{\pi} g_{5/2}(Z) = \frac{1}{\lambda^5} g_{5/2}(Z)$$

$$\rho = \frac{1}{V} \frac{Z}{1-Z} + \frac{1}{\lambda^5} g_{5/2}(Z)$$

• Nota $g_{5/2}(1) < \infty$, quindi c'è condensazione

$$(3) z \frac{\partial}{\partial z} \langle N \rangle = \frac{z}{1-z} + \frac{z^2}{(1-z)^2} + \frac{V}{\lambda^3} g_{3/2}(z) = \langle N^2 \rangle - \langle N \rangle^2$$

$$\rho > \rho_c(\pi) = \frac{g_{5/2}(1)}{\lambda^3(\pi)}$$

$$(4) V \rightarrow \infty \quad z \approx 1 - \frac{1}{N_0}$$

$$\frac{z}{1-z} \approx \frac{1 - \frac{1}{N_0}}{1 - 1 + \frac{1}{N_0}} = N_0 - 1 \approx N_0, \quad \langle N \rangle = N_0 + \frac{V}{\lambda^3} g_{3/2}(1) = N_0 + V \rho_c = N_0 + N_c$$

$$\langle N^2 \rangle - \langle N \rangle^2 = -N_0 + N_0^2 + \frac{V}{\lambda^3} g_{3/2}(1)$$

$$\frac{\langle N^2 \rangle - \langle N \rangle^2}{\langle N \rangle^2} = \frac{N_0^2 - N_0 + \frac{V}{\lambda^3} g_{3/2}(1)}{(N_0 + \frac{V}{\lambda^3} g_{3/2}(1))^2} \approx \frac{N_0^2}{\langle N \rangle^2} = \frac{\rho_0^2}{\rho^2}$$

$$\sqrt{\frac{\langle N^2 \rangle - \langle N \rangle^2}{\langle N \rangle^2}} \approx \frac{\rho_0}{\rho} \sim 1$$

Le fluttuazioni non sono trascurabili, sono macroscopiche!